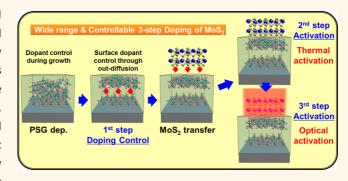
Wide-Range Controllable n-Doping of Molybdenum Disulfide (MoS₂) through Thermal and Optical Activation

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ABSTRACT Despite growing interest in doping two-dimensional (2D) transition metal dichalcogenides (TMDs) for future layered semiconductor devices, controllability is currently limited to only heavy doping (degenerate regime). This causes 2D materials to act as metallic layers, and an ion implantation technique with precise doping controllability is not available for these materials (e.g., MoS₂, MoSe₂, WSe₂, graphene). Since adjustment of the electrical and optical properties of 2D materials is possible within a light (nondegenerate) doping regime, a wide-range doping capability including nondegenerate and degenerate regimes is a critical aspect



of the design and fabrication of 2D TMD-based electronic and optoelectronic devices. Here, we demonstrate a wide-range controllable n-doping method on a 2D TMD material (exfoliated trilayer and bulk MoS₂) with the assistance of a phosphorus silicate glass (PSG) insulating layer, which has the broadest doping range among the results reported to date (between 3.6×10^{10} and 8.3×10^{12} cm⁻²) and is also applicable to other 2D semiconductors. This is achieved through (1) a three-step process consisting of, first, dopant out-diffusion between 700 and 900 °C, second, thermal activation at 500 °C, and, third, optical activation above 5 μ W steps and (2) weight percentage adjustment of P atoms in PSG (2 and 5 wt %). We anticipate our widely controllable n-doping method to be a starting point for the successful integration of future layered semiconductor devices.

KEYWORDS: MoS₂ · wide-range · controllable doping · thermal activation · optical activation

olybdenum disulfide (MoS₂), a transition metal dichalcogenide (TMD) with layered structure, has recently been considered a promising candidate for next-generation flexible electronic and optoelectronic devices because of its superior electrical, 1-4 optical, 5-7 and mechanical properties.⁸ Thickness scalability down to a monolayer and van der Waals expitaxial structure without surface dangling bonds (consequently, native oxides) make MoS₂-based thin film transistors (TFTs) immune to the short channel effect (SCE) and provides very high field-effect mobility (\sim 200 cm²/(V s), which is comparable to the universal mobility of Si), respectively.^{1,9} In addition, an excellent photodetector with

a wide spectral range from ultraviolet (UV) to near-infrared (IR) is achievable using MoS₂, since its energy band gap varies between 1.2 eV (bulk) and 1.8 eV (monolayer), depending on layer thickness.^{6,7,9–12} It also has potential for use in chemical sensors, ^{13,14} memory, ^{15,16} and valleytronic devices. ^{4,5,17} However, one of the critical issues that hinder the successful integration of such MoS₂ electronic and optoelectronic devices is the lack of a reliable and controllable doping method. Such a component is essential for inducing a shift in the Fermi level, which subsequently enables wide modulations of its electrical and optical properties.^{18,19}

Fang et al. reported a potassium-based degenerate n-doping process via a surface

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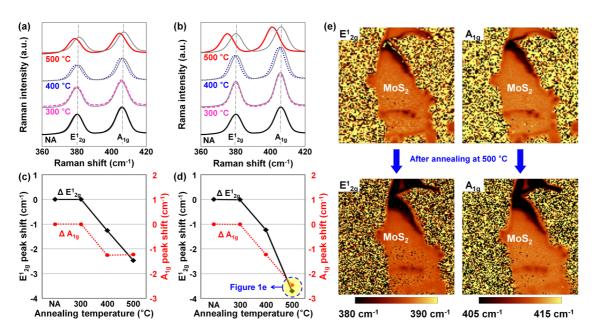


Figure 1. Raman spectrum analysis of MoS_2 on SiO_2 and PSG. Raman spectra of MoS_2 on (a) SiO_2 and (b) PSG, which were nonannealed (NA) and annealed at different temperatures (300, 400, and 500 °C). E^1_{2g} and A_{1g} peak shift data as a function of annealing temperature, which were extracted from the Raman spectra of (c) MoS_2/SiO_2 and (d) MoS_2/PSG samples. (e) Raman mapping images in E^1_{2g} and A_{1g} peaks for the MoS_2 flake before/after the n-doping process through 500 °C annealing.

charge transfer mechanism and then used it to form highly n-doped source/drain (S/D) regions in TMD-based TFTs.²⁰ Sreeprasad et al. observed n- and p-doping phenomena on MoS₂ through different functionalization methods of Au nanoparticles.²¹ In addition, n- and p-doping techniques based on polymers such as polyethylenimine (PEI) and functional selfassembled monolayers (SAMs) with different dipole moments were reported by Du et al.²² and Li et al.²³ Chen et al. also presented a p-doping method on MoS₂ by using selected-area plasma treatment with fluorine (F) or oxygen (O).²⁴ Although Fang et al. controlled the doping level by adjusting the exposure time to potassium, the range of doping was limited within the heavily doped (degenerate) regime, where MoS₂ works as a near-metallic layer. Since adjusting the electrical and optical properties of MoS₂ is possible within a nondegenerate doping regime, wide-range doping capability including nondegenerate and degenerate regimes is a critical point in the design and fabrication of MoS₂-based electronic and optoelectronic devices. However, it is very challenging to achieve nondegenerate doping on 2D semiconductors and secure its widerange doping controllability because a method such as ion implantation cannot be applied to the doping process. Recently, Lin et al. reported nondegenerate n-doping of MoS₂ using cesium carbonate (Cs₂CO₃).²⁵ In this work, we first demonstrate a wide-range controllable n-doping method for MoS₂ on a phosphorus silicate glass (PSG) insulating layer, which (1) activates the doping phenomenon via a thermal/optical process and (2) adjusts the doping level of MoS₂ by controlling the thermal/optical process conditions or weight percentage of P atoms during the in situ doped PSG

growth step. This controllable n-doping process is also demonstrated on exfoliated trilayer and bulk MoS_2 films. In addition, the proposed doping method is investigated in detail through Raman spectroscopy, X-ray photoelectron spectroscopy (XPS), electrical measurements (I_D – V_G and I_D – V_D), and atomic force microscopy (AFM).

RESULTS AND DISCUSSION

Characterization of n-Doped MoS2. First, Raman spectroscopy analysis was performed on the MoS₂/SiO₂ and MoS₂/PSG samples to investigate the effects of (1) the thermal annealing process on MoS₂ and (2) P atoms (these atoms exist in the form of phosphorus pentoxide, a chemical compound with the molecular formula P₂O₅) in the PSG layer on MoS₂ respectively, during second-step annealing. Figure 1a and b show the Raman spectra measured on MoS₂/SiO₂ and MoS₂/ PSG samples with or without annealing. Here, two conventional peaks (E¹_{2q} and A_{1q}) indicating the inplane and out-of-plane vibrations for bulk MoS₂ were respectively observed at \sim 380 and \sim 406 cm⁻¹. Then, the peak position shifts in E_{2q}^1 and A_{1q} after the annealing process were extracted and plotted as a function of annealing temperature, as shown in Figure 1c and d. No shift in E_{2q}^1 and A_{1q} peaks was observed for MoS₂/SiO₂ and MoS₂/PSG samples after a 300 °C annealing step compared to their initial peak positions. Although the peaks were slightly red-shifted by \sim 1.2 cm⁻¹ for the 400 °C annealed MoS₂/PSG sample, this seems to originate from the thermal annealing effect on MoS₂ because it was also observed in the 400 °C annealed MoS₂/SiO₂ sample. This redshift phenomenon is mainly attributed to the changes

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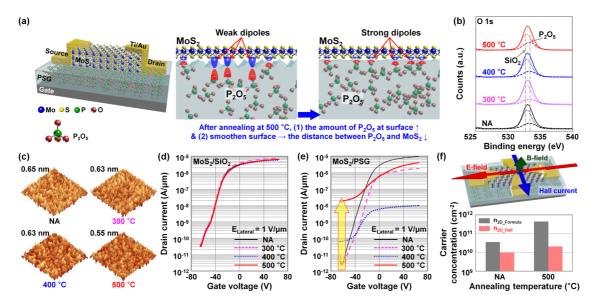


Figure 2. One-step n-doping through thermal activation. (a) Schematics of back-gated transistors fabricated on MoS₂/PSG, explaining the reaction between MoS₂ and P₂O₅ in the NA and 500 °C annealed samples. (b) XPS (O 1s peak intensity as a function of binding energy) and (c) AFM surface analysis of MoS2 films on PSG substrates, which were nonannealed and annealed at 300, 400, and 500 $^{\circ}$ C. $I_D - V_G$ characteristics of MoS₂ transistors fabricated on (d) SiO₂ and (e) PSG substrates, which were nonannealed and annealed at 300, 400, and 500 °C. Here, the ID was normalized by channel width (W). (f) Comparison between the carrier concentrations obtained from the formula and the Hall-effect measurement.

in E_{2q}^1 and A_{1q} mode vibrations caused by thermal expansion of MoS₂ after 400 °C annealing.²⁶ Similar red-shift phenomenon in the E¹_{2q} peak was previously reported by Tongay et al. after 450 °C annealing for 40 min, and this was possibly attributed to the desorption of contaminants or the stress-releasing effect on MoS₂.²⁷ However, for higher temperature annealing (500 °C), the peaks in the MoS₂/PSG sample are pushed further to the left ($\Delta E_{2q}^1 = \sim 3.7 \text{ cm}^{-1}$ and $\Delta A_{1q} =$ \sim 2.46 cm $^{-1}$) when compared with the peak shifts $(\Delta E_{2q}^{1} = \sim 2.47 \text{ cm}^{-1} \text{ and } \Delta A_{1q} = \sim 1.23 \text{ cm}^{-1}) \text{ in a}$ MoS₂/SiO₂ sample. P₂O₅ molecules in the surface region of the PSG layer are thought to electrically affect MoS₂ after 500 °C annealing, which consequently provides more electrons to MoS₂. Sreeprasad et al. previously reported the red-shift of E_{2q}^1 and A_{1q} peaks after the n-doping process on MoS2, which induced functionalization of Au nanoparticles on MoS₂ through instantaneous reaction using microwave.²¹ In addition, Li et al. observed the red-shift of the A_{1a} peak by 1.6 cm⁻¹ in MoS₂ flakes on a 3-(trimethoxysilyl)-1propanamine (APTMS)-treated substrate and confirmed the n-doping phenomenon through electrical measurements.²³ The distance between E_{2q}^1 and A_{1q} peaks obtained in all of the samples in the present study was larger than 24 cm⁻¹, since measurement was done on thick MoS₂ flakes, considered bulk MoS₂. Figure 1e shows the Raman mapping images corresponding to E_{2q}^1 and A_{1q} peaks before/after the n-doping process through 500 °C annealing. Darker MoS₂ images were observed after the annealing process, indicating that the MoS₂ flake was uniformly n-doped by the P atoms in PSG. Although there are color

differences in each Raman mapping image, which seem to be caused by thickness variation in the exfoliated MoS2 flake, this does not indicate nonuniformity of the n-doping concentration, because E¹_{2q} and A_{1q} peak shift values are constant at all positions despite thickness variation.

Figure 2a shows the schematics of the back-gated transistor fabricated on MoS₂ and explains the n-doping mechanism. Since the P₂O₅ molecule in the PSG layer is asymmetric in structure with nonzero dipole momentum (having a positive pole in P atoms and a negative pole in O atoms), the n-doping effect is thought to originate from another dipole induction at the MoS₂/PSG interface. The O atoms with a negative pole on the PSG surface seem to attract holes from the MoS₂ layer and hold them at the interface region, thereby n-doping the MoS₂ layer. In addition, the increased amount of P2O5 molecules on the PSG surface and the annealing-related surface smoothing are predicted to strengthen the dipole moment and, consequently, the degree of n-doping. As shown in Figure 2b, the XPS spectra of O 1s peaks, the intensity of the P₂O₅ signal (at 533.4 eV) increased slightly as the annealing temperature rose to 500 °C, consequently indicating the larger amount of P2O5 in the surface region. Figure 2c presents AFM images and surface roughness on the PSG samples, which are nonannealed or annealed between 300 and 500 °C. Although it was confirmed that P2O5 molecules were outdiffused from PSG film via annealing, the rootmean-square (RMS) roughness of the annealed PSG film was reduced from 0.65 nm to 0.55 nm as the annealing temperature increased to 500 °C due to the

high reflowing property of the PSG film. As a result, a shorter distance between P₂O₅ (particularly the negatively charged O atoms) and MoS2 is expected in the 500 °C annealed sample when compared to other samples annealed at lower temperatures. In order to reconfirm the n-doping phenomenon on MoS₂ verified by Raman spectroscopy analysis in Figure 1, the $I_D - V_G$ characteristics of MoS₂ transistors were measured and compared in Figures 2d and e, where the lateral electric field ($E_{Lateral}$) between the source and drain electrodes was fixed at 1 V/ μ m. Although shifts in E_{2q}^{1} and A_{1q} peaks were observed on MoS₂/SiO₂ samples that were annealed at 400 and 500 °C, no changes were observed in the $I_D - V_G$ curves of Figure 2d, indicating that the peak shifts did not originate from the n-doping effect. However, for the case of MoS₂/PSG samples, a rising trend in off-current (2 \times 10⁻¹² A/ μ m \rightarrow 6 \times 10⁻¹¹ A/ μ m \rightarrow 2 \times 10⁻⁸ A/ μ m at V_{GS} = -55 V) and a negative shift in threshold voltage ($-25 \text{ V} \rightarrow -40 \text{ V} \rightarrow -42 \text{ V}$) were confirmed as the annealing temperature increased from 300 °C to 500 °C. Additional electrons are thought to be supplied to the MoS₂ after annealing above 400 °C because the n-doping phenomenon associated with the formation of dipoles eventually influences the tunneling of electron carriers from the source metal to MoS₂. The MoS₂ energy band is downshifted due to n-doping, which increases the electric field at the source-MoS₂ junction, thereby enhancing the tunneling of electrons and causing a negative V_{TH} shift. The increase in off-current can be explained by the interfacial charge layer caused by the dipoles, which consists of hole carriers. Although the injection of hole carriers from drain to MoS₂ is not significant even under negative V_{GS} , this conductive interfacial hole layer is expected to induce the transport of hole carriers and a subsequent increase in off-currents. This n-doping phenomenon is also present in the $I_D - V_D$ characteristic curves (Supporting Information Figure S1) obtained for the nonannealed (NA) and the 500 °C annealed MoS₂/PSG samples. For the 500 °C annealed MoS_2/PSG device, the drain current (I_D) did not saturate as drain voltage (V_{DS}) increased, because a higher n-type carrier concentration is formed in the nonaccumulated region (between the drain electrode edge and the channel edge) after the pinch-off regime ($V_{DS} > V_{GS} - V_{TH}$). The 2D sheet doping concentration (n_{2D}) extracted from $n_{\rm 2D} = I_{\rm D} L/qW\mu V_{\rm DS}$ is also consistent with the rising trend in off-current as a function of second-step annealing temperature, showing a significant increase from 3.6 \times $10^{10}\,\text{cm}^{-2}$ to $4.5\times10^{11}\,\text{cm}^{-2}$ (Supporting Information Figure S2a). In the doping concentration equation, L and W refer to the length and width of the MoS₂ channel, respectively, q is the electron charge, μ is the field-effect mobility at $V_{GS} = 0 \text{ V}$, I_D is the drain current at $V_{GS} = 0 \text{ V}$, and $V_{\rm DS}$ is the drain—source voltage. The obtained $n_{\rm 2D}$ values are much lower than the concentrations (1.0 \times 10^{13} and 7.2×10^{12} cm⁻²) achieved by the potassium²⁰

and APTMS²³ doping processes, which were performed on three- and one-layer MoS2, respectively. In addition, on-currents in MoS₂/PSG devices were dramatically changed as a function of annealing temperature, unlike MoS₂/SiO₂ samples, as shown in Figure 2e. The oncurrent was first reduced when the annealing temperature was increased up to 400 °C, but was increased again after going through 500 °C annealing. Compared to the control and 300 °C annealed samples, more P₂O₅ molecules were expected to gather at the surface region of the PGS layer and cause a slightly stronger reaction with MoS₂. The negative poles of O atoms in P₂O₅ molecules influence the MoS₂ channel region and consequently seem to cause more severe interfacial carrier scattering, eventually decreasing on-current level and field-effect mobility in the n-channel MoS₂/PSG transistor. According to the maximum field-effect mobility extracted at $V_{DS} = 5 \text{ V}$ in Supporting Information Figure S2b, the initial value (10.5 cm² V⁻¹ s⁻¹) of the control device was severely decreased down to 0.01 cm² V⁻¹ s⁻¹ when the second-step annealing temperature increased to 400 °C. However, on-current was recovered to a similar level to the control sample after 500 °C annealing, and we also confirmed that the mobility was increased up to 3.9 cm 2 V $^{-1}$ s $^{-1}$ in Figure S2b. It is thought that MoS $_2$ reacts more strongly with P2O5 molecules in PSG at this temperature and the large amount of electrons provided to MoS₂ consequently increase the on-current. As shown in Figure 2f, we also note that the carrier concentrations $(3.6 \times 10^{10} \text{ cm}^{-2} \text{ for NA-NA and } 4.5 \times 10^{11} \text{ cm}^{-2} \text{ for NA-NA})$ NA-500 °C) obtained from the formula were slightly overestimated when compared to those $(9.8 \times 10^9 \text{ cm}^{-2})$ for NA-NA and 2.1 \times 10¹⁰ cm⁻² for NA-500 °C) obtained from Hall-effect measurement.

Two-Step n-Doping Process through Thermal Activation. To adjust the amount of P2O5 molecules in the surface region through out-diffusion, additional annealing steps (first annealing step) between 700 and 900 °C were performed on the PSG films in a N₂ ambient, as shown in Figure 3a. Comparing the XPS spectra for O 1s peaks between the as-deposited PSG film and the PSG films mentioned above (Figure 3b), the intensity of the P₂O₅ signal (at 533.4 eV) increased after annealing at higher temperatures. In addition, the atomic percentage of the P 2p XPS peak (at 134.58 eV) gradually increased as a function of annealing temperature (0.7% at NA, 0.77% at 700 °C, 1.74% at 800 °C, and 2.02% at 900 °C in Supporting Information Figure 3a). These indirectly indicate a higher surface concentration of P atoms and, consequently, a larger number of P₂O₅ molecules after higher temperature first-step annealing. Since it was confirmed that 500 °C annealing is the minimum requirement to activate the surface doping phenomenon, all second-step annealing after MoS₂ exfoliation was done at 500 °C. Figure 3c shows the E_{2a}^{1} and A_{1a} peak shifts in MoS₂ after the first and second annealing steps, which were extracted from

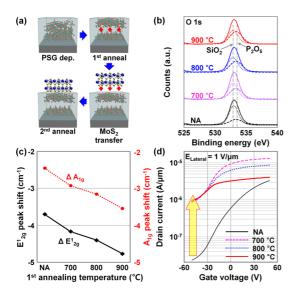


Figure 3. Two-step n-doping through thermal activation. (a) Schematic diagram showing the two-step ${\rm MoS}_2$ doping process. (b) XPS surface analysis (O 1s peak intensity as a function of binding energy) of PSG films, which were non-annealed and annealed at 700, 800, and 900 °C (first-step annealing). (c) Extracted Raman peak shifts in ${\rm E}^1_{2g}$ and ${\rm A}_{1g}$ after the two-step ${\rm MoS}_2$ doping process as a function of the first-step annealing temperature. (d) $I_{\rm D}$ – $V_{\rm G}$ characteristics of ${\rm MoS}_2$ transistors fabricated without/with different first-step annealing at 700, 800, and 900 °C. Here, the $I_{\rm D}$ was normalized by channel width (W).

the MoS₂/PSG Raman spectra in Supporting Information Figure S3b. E_{2q}^1 and A_{1q} peaks in the 700–500 °C (700 °C first-step and 500 °C second-step annealing) annealed MoS₂/PSG sample were more red-shifted (by \sim 4.17 and \sim 2.94 cm⁻¹) compared to the peak shifts (\sim 3.7 and \sim 2.46 cm $^{-1}$) of the NA-500 °C (no first-step and 500 °C second-step annealing) annealed MoS₂/ PSG sample. In addition, the peaks were red-shifted more and more as the first-step annealing temperature increased from 700 °C to 900 °C (\sim 4.77 cm⁻¹ in $E_{2\alpha}^{1}$ and \sim 3.55 cm⁻¹ in A_{1a}). This indicates higher n-doping of MoS₂ due to the increase in P₂O₅ molecules in the surface region and the reduced distance between P₂O₅ and MoS₂ after first-step annealing. This doping controllability was confirmed once again by electrical measurement $(I_D - V_G)$ of back-gated transistors fabricated on MoS₂/PSG samples. Due to the higher n-doping level achieved via first-step annealing, off-currents per width (at $V_{GS} = -50$ V and $E_{Lateral} = 1$ V/ μ m) of these transistors were 8.4×10^{-7} , 8.9×10^{-7} , and 9.4×10^{-7} 10^{-7} A/ μ m, which were much higher than that of the previous transistor formed on the MoS₂/PSG sample without first-step annealing (2.2 \times 10⁻⁸ A/ μ m at V_{GS} = -50 V), indicating a higher n-doping level in the sample exposed to a higher process temperature. The extracted n_{2D} of the 900–500 °C sample was about $1.0 \times 10^{12} \ \text{cm}^{-2}$, which is slightly higher than that of the NA-500 °C sample (4.5 \times 10¹¹ cm $^{-2}$). Meanwhile, a decreasing trend in the on-current with temperature increase was observed due to the higher concentration

of P atoms in the interfacial region, inducing extensive carrier scattering. On the basis of AFM analysis performed on the PSG samples (Supporting Information Figure S3c), the RMS surface roughness was reduced as the process temperature rose, owing to the high reflowing property of the PSG film. Through this, it can be predicted that interfacial scattering was mainly dominated by the Columbic scattering effect due to P atoms at the interface, not by physical scattering due to the surface roughness of the PSG. The n-doping levels that were achieved using a one- or two-step doping process were confirmed once again with XPS analysis in Supporting Information Figure S4. These upshifts observed in the 900-500 °C annealed sample (compared to NA-NA) were \sim 0.2 eV (Mo 3d_{5/2}), \sim 0.3 eV (Mo 3d_{3/2}), \sim 0.2 eV (S 2p_{1/2}), and \sim 0.1 eV (S 2p_{3/2}), directly indicating successful n-doping on MoS2 because the Fermi level shifted toward the conduction band edge because n-doping results in a lower binding energy. Although this shift is coincident with previously reported XPS data on potassium-doped MoS₂, the degree of shift obtained by this work is smaller than that in the previous report (\sim 0.7 eV in the Mo 3d peak and \sim 0.6 eV in the S 2p peak), which was degenerately doped by potassium.²⁰ This also indicates that the n-doping levels of our one- or two-step thermal process on PSG are in the nondegenerate regime. In addition, the achieved nondegenerate n-doping phenomenon was stable even after exposure in air for 4 days, based on the I_D-V_G characteristic of the 900-500 °C annealed device sample (Supporting Information Figure S5). As seen in the control sample (NA-NA), a small reduction in drain current also occurred in the doped device (900-500 °C).

As shown in Figure 4a, another PSG film with a lower P weight percentage (2 wt %) than the previous film (5 wt %) was used to adjust the concentration of P atoms on the surface of the PSG film and eventually the n-doping level. In this doping control experiment, firstand second-step annealing process temperatures were fixed at 900 and 500 °C, respectively, to clearly observe differences in the level of n-doping. According to XPS analysis performed on 5 and 2 wt % samples in Figure 4b and Supporting Information Figure S6a, where the intensity of P₂O₅ and P in the 2 wt % sample was lower than that in 5 wt %, a slightly lower surface concentration of P atoms (thereby, P2O5 molecules) in the 2 wt % sample was confirmed after first-step annealing at 900 °C. As a result, after performing second-step annealing at 500 °C to activate the n-doping phenomenon, smaller peak shifts ($\Delta E_{2q}^1 = \sim 2.96 \text{ cm}^{-1}$ and $\Delta A_{1g} = \sim 1.72 \text{ cm}^{-1}$) were observed in the 2 wt % sample than seen in the 5 wt % sample $(\Delta E_{2q}^1 =$ \sim 3.7 cm⁻¹ and $\Delta A_{1q} = \sim$ 2.46 cm⁻¹) in Figure 4c. Since the peak shifts in the 2 wt % sample were slightly larger than those ($\Delta E_{2q}^1 = \sim 2.47 \text{ cm}^{-1}$ and $\Delta A_{1q} = \sim 1.23 \text{ cm}^{-1}$) of the 500 °C annealed MoS₂/SiO₂ sample, MoS₂ on

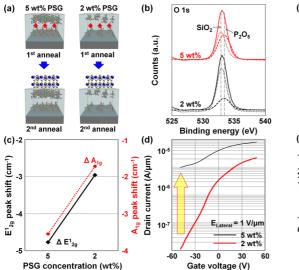


Figure 4. Two-step n-doping through thermal activation on different PSG substrates. (a) Schematic diagram showing the MoS_2 doping process performed on PSG films with different P weight percentages (5 and 2 wt %). (b) XPS surface analysis (O 1s peak intensity as a function of binding energy) of the PSG films. (c) Extracted Raman peak shifts in E^1_{2g} and A_{1g} as a function of P weight percentage. (d) $I_D - V_G$ characteristics of MoS_2 transistors fabricated on 5 and 2 wt % PSG films. Here, the I_D was normalized by channel width (W).

2 wt % PSG film is expected to be successfully n-doped at a much lower doping level. Back-gated MoS2 transistors fabricated on the 5 and 2 wt % PSG films also demonstrate a significant difference in off-currents at $V_{\rm GS} = -50$ V (\sim 1.2 \times 10⁻⁷ A/ μ m in 5 wt % and \sim 4.5 \times 10^{-10} A/ μ m in 2 wt %) in Figure 4d, which provide evidence of the lower n-doping level of MoS₂ on 2 wt % PSG film. Although a decreasing trend was previously observed in the on-current due to the interfacial scattering effect in the high doping level regime (Figure 2), in this study, higher on-current was obtained in the 5 wt % sample with a higher surface concentration of P2O5. The dramatically increased number of electrons transferred from source metal to MoS₂ as the P₂O₅ concentration of PSG rose from 2 to 5 wt % has an ascendancy over the scattering effect.

Three-Step n-Doping Process via Thermal and Optical Activation. Additional laser exposure was performed on the 900-500 °C sample to enhance the activation of the n-doping phenomenon. Here, we also note that the second-step thermal anneal above 500 °C cannot be used to achieve higher n-doping concentration because it decomposes MoS₂ flakes even in H₂ ambient. Figure 5a shows a schematic diagram briefly explaining the optical n-doping process. When considering the energy band gap (\sim 1.2 eV) of thick MoS₂ flakes, higher photon energy incident lasers are predicted to cause extra energy loss ($hv - E_g$) to reach thermal equilibrium, consequently providing heat energy to the MoS₂/PSG interface and enhancing n-doping activation. Because MoS₂ exhibits different light absorption properties as the wavelength varies, three laser sources

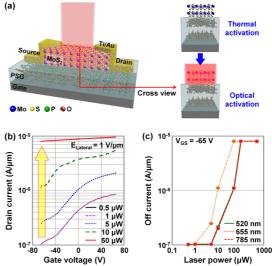


Figure 5. Three-step n-doping via thermal and optical activation. (a) Schematic diagrams explaining the optical MoS₂ n-doping process performed on the 900–500 °C annealed MoS₂/PSG sample. (b) $I_{\rm D}-V_{\rm G}$ characteristics of MoS₂ transistors that were exposed to a 655 nm wavelength laser with power varying from 0.5 to 1000 μ W. (c) Extracted off-currents (at $V_{\rm GS}=-65$ V) as a function of laser power after exposing the transistors to lasers at three different wavelengths ($\lambda=520$, 655, and 785 nm).

with wavelengths of 520, 655, and 785 nm were used in this experiment. Figure 5b shows the $I_D - V_G$ characteristics of back-gated MoS₂ transistors that were exposed to the 655 nm wavelength laser with power varying from 0.5 to 1000 μ W. Although no change in the $I_D - V_G$ curves was observed below 1 μ W, the current level was increased about 70-fold when the power rose from $5 \,\mu\text{W}$ to $50 \,\mu\text{W}$, indicating that the additional n-doping effect occurred above 5 μW for the 655 nm laser. In particular, the n-doping concentration increased from 1.0 \times 10¹² cm⁻² (900–500 °C sample) to 8.3 \times 10^{12} cm^{-2} after exposure to the 10 μW laser, and the doped MoS₂ eventually exhibited metallic properties at 50 μ W. This means that the n-doping level of MoS₂ can be successfully controlled in the degenerate doping regime by adjusting the laser power. We then plotted the off-current values (at $V_{\rm GS} = -65$ V) extracted from the $I_D - V_G$ curves as a function of the laser power in Figure 5c after exposing the transistors to lasers with three different wavelengths. In the cases where 520 nm (shorter) and 785 nm (longer) wavelengths were used, the additional laser-induced n-doping phenomenon happened at a slightly higher power point of 10 μ W, owing to the lower light absorption at the wavelengths than that at 655 nm. This indicates that laser wavelength adjustment, in addition to laser power adjustment, can also be a useful technique for n-doping control. For reference, we also provide $I_D - V_G$ characteristic data for MoS₂ transistors on SiO₂ and PSG for a number of laser illumination cycles (on/off) in Supporting Information Figure S7. Although the drain current increased by laser exposure returned to the

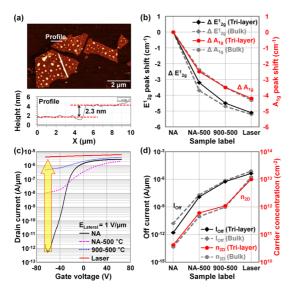


Figure 6. Three-step thermal and optical n-doping on trilayer and bulk MoS $_2$. (a) AFM images and height ($\sim\!2.3$ nm) of the trilayer MoS $_2$ flake. (b) Extracted Raman peak shifts in E $^1_{2g}$ and A $_{1g}$ obtained on trilayer and bulk MoS $_2$ samples (NA–NA, NA–500 °C, 900–500 °C, 900–500 °C + laser). (c) $I_D\!-\!V_G$ characteristics of MoS $_2$ transistors fabricated on trilayer MoS $_2$ samples. Here, I_D was also normalized by channel width (W). (d) Off-currents and carrier concentrations extracted from the $I_D\!-\!V_G$ curves of trilayer and bulk MoS $_2$ samples.

dark current level when the laser was turned off in devices on SiO₂, the current of the device on PSG was not altered before or after laser illumination.

Three-Step Thermal and Optical Doping on Trilayer and Bulk MoS₂. The thermal and optical doping process was performed again on trilayer MoS₂ films to determine whether a thickness effect influences n-doping. As shown in Figure 6a, the thickness of MoS₂ flakes was reduced down to \sim 2.3 nm (trilayer) from \sim 30 nm (bulk). ΔE_{2q}^1 and ΔA_{1q} values extracted from Raman peaks before/after the doping process are plotted and compared with those of bulk MoS₂ in Figure 6b. The ΔE_{2q}^1 and ΔA_{1q} values of trilayer MoS₂ respectively decreased to -5.1 cm^{-1} and -4.2 cm^{-1} after first-step annealing (at 900 °C) and laser treatment (λ = 655 nm and $P = 1000 \mu W$); trilayer MoS₂ showed similar behavior to bulk MoS₂ (gray dotted lines). In addition, the MoS₂ transistors fabricated under four different doping conditions (NA-NA, NA-500 °C, 900-500 °C, and 900-500 °C + laser exposure) demonstrated significantly different off-currents (\sim 1.67 \times 10⁻¹², \sim 1.25 \times 10^{-8} , $\sim 5.75 \times 10^{-7}$, and $\sim 4.41 \times 10^{-6}$ A/ μ m, respectively, at $V_{\rm GS} = -55$ V) in Figure 6c and d. However, the increasing trend in the off-current (black solid line) is similar to that observed in the bulk MoS₂ transistors (gray dotted line). In particular, the $n_{\rm 2D}$ values of NA-500 °C, 900-500 °C, and 900-500 °C + laser samples (trilayer MoS₂ transistors) are 6.1×10^{11} , 1.1×10^{11} 10^{12} , and 1.0×10^{13} cm⁻², which are comparable to those obtained for bulk MoS₂ transistors (4.5 \times 10¹¹, 1.0×10^{12} , and 1.2×10^{13} cm⁻²). These results indicate that the doping range is very broad from nondegenerate to degenerate and also the thickness of MoS₂ does not have an influence on the n-doping phenomenon. In previous cases of degenerate n-doping of MoS₂ via potassium²⁰- and APTMS²³-doping processes, the n-doping concentrations were 1.0 \times 10^{13} and 7.2 \times 10¹² cm⁻², which were very similar to that of our 900-500 °C + laser sample. The doping concentration of nondegenerate n-doped MoS₂ by cesium carbonate (Cs₂CO₃) was between 1×10^{11} and 3.5×10^{11} cm^{-2,25} which is also close to that of the NA-500 C sample.

CONCLUSIONS

In summary, a widely controllable n-doping process for MoS₂ on a PSG insulating substrate was demonstrated by adjusting the thermal and optical process conditions (T = 300-500 °C for the thermal process and $P = 0.5 - 1000 \,\mu\text{W}$ for the optical process) or the concentration of P atoms in the PSG (5 and 2 wt %, thereby affecting the concentration of P₂O₅ molecules). It was confirmed that 500 °C annealing and laser exposure at 5 μ W (λ = 655 nm) to 10 μ W (λ = 520 and 785 nm) were the minimum requirements to activate the surface doping phenomenon. Through this, MoS2 was uniformly n-doped by P atoms in PSG, which was confirmed through Raman spectroscopy and XPS analysis. The controlled doping levels were between 3.6×10^{10} and $8.3 \times 10^{12} \text{ cm}^{-2}$ before formation of a metallic layer, where MoS₂-based electronic and optoelectronic devices can be properly designed. The doping range is broadest among the results reported up to now, and also, this method is applicable to other 2D semiconductors (e.g., MoS₂, MoSe₂, WS₂, WSe₂, graphene). This was achieved by (1) a three-step process consisting of (i) dopant outdiffusion between 700 and 900 °C, (ii) thermal activation at 500 °C, and (iii) optical activation above 5 μ W steps and (2) weight percentage adjustment of P atoms (thereby, P₂O₅ molecules) in PSG. This controllable n-doping effect was also demonstrated on exfoliated trilayer and bulk MoS₂, where it was confirmed that the thickness of MoS₂ does not influence the n-doping phenomenon. We expect our wide-range controllable n-doping method to be a starting point for the successful integration of future layered semiconductor devices.

EXPERIMETAL METHODS

Preparation of PSG Films and Surface Doping Control. First, 300 nm thick PSG layers with a 5 or 2 wt % P doping concentration were deposited on a heavily doped p-type Si substrate (resistivity

<0.005 Ω -m) at 480 °C in a low-pressure chemical vapor deposition (LPCVD) system. In order to control the amount of P atoms in the surface region, we performed a first-step annealing process in N₂ ambient from 700 to 900 °C.

Thermal and Optical Doping of MoS₂ on PSG Films. After transferring MoS₂ films on the annealed and nonannealed PSG samples using adhesive tape, an acetone wash was done on the samples for 1 h to remove tape residue. Then, a second-step thermal activation process was carried out in N₂/H₂ ambient from 300 to 500 °C, or optical activation (refer to Supporting Information Figure S7) was performed with laser sources at wavelengths of 520, 655, and 785 nm, in order to cause a reaction between the P atoms in the surface region of the PSG and MoS₂, eventually doping MoS₂. The doping level of MoS₂ was controlled by adjusting the second-step thermal/optical process conditions as well as regulating the weight percentage of P atoms during the *in situ* doped PSG growth step.

Characterization of Doped MoS₂ Films. The prepared MoS₂/PSG/Si samples were investigated and compared with control samples (MoS₂/SiO₂/Si) by PL/Raman spectroscopy (Alpha300 M+, WITec), XPS (ESCA200, VG Microtech Inc.), and AFM (SPA-300, Seiko Instrument) measurements. Raman spectroscopy with an excitation wavelength of 532 nm was used; its laser beam size was 0.7–0.9 μ m, and the instrumental spectral resolution was below 0.9 cm $^{-1}$. An integration time of 5 s and spectrometry at 1800 grooves/mm were employed for the test. A Mg K α twinanode source and 0° of X-ray incident angle were used for XPS measurement.

Fabrication and Electrical Measurements of Undoped and Doped MoS₂ Devices. For the fabrication of back-gated MoS₂ devices (transistors and photodetectors), source/drain electrode regions were patterned on the MoS₂/PSG/Si and MoS₂/SiO₂/Si samples using optical lithography to achieve a 1 μ m (or 5 μ m) channel length and 1 μ m (or 5 μ m) channel width, followed by Ti (10 nm) and Au (50 nm) deposition via e-beam evaporator. For the fabrication of trilayer MoS₂ transistors, an e-beam lithography process was used instead of optical lithography. The fabricated devices were analyzed through electrical measurements $(I_D - V_G)$ and $I_D - V_D$) by an HP 4155B semiconductor parameter analyzer, and we subsequently calculated the carrier concentration and field-effect mobility in the ${\mbox{MoS}}_2$ accumulated channel using the measured data. Here, all drain currents (I_D) were normalized by channel width (W). The n_{2D} was extracted from n_{2D} = $I_DL/qW\mu V_{DS}$, where q is the electron charge, μ is the field-effect mobility at $V_{GS} = 0 \text{ V}$, I_D is the drain current at $V_{GS} = 0 \text{ V}$, and V_{DS} is the drain—source voltage. In addition, Hall-effect measurement was also performed to confirm the extracted n_{2D} values. For field-effect mobility extraction, the equation $\mu_{FE} = L/(WV_{DS}C_{OX}) \times$ $(\partial I_D/\partial V_{GS})$ was used, where L and W are the length and channel width, and $C_{\rm OX}$ is $\varepsilon_{\rm OX} \times \varepsilon_{\rm 0}/t_{\rm OX}$, which is the gate oxide capacitance per unit area.

Conflict of Interest: The authors declare no competing financial interest.

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Supporting Information Available: $I_{\rm D}$ – $V_{\rm D}$ characteristics of nonannealed and 500 °C annealed devices. Extraction of electrical parameters (carrier concentration and field-effect mobility) as a function of second-step annealing temperature (no first-step annealing). XPS, Raman spectrum, and AFM analysis of PSG films annealed at different temperatures. XPS analysis of MoS $_2$ /PSG samples (NA, NA–500 °C, 900–500 °C). Air-stability analysis of undoped and doped MoS $_2$ devices. XPS and Raman spectrum analysis of PSG films with different P concentration. Experimental setup for the laser annealing process and $I_{\rm D}$ – $V_{\rm G}$ characteristics of MoS $_2$ transistors on SiO $_2$ and PSG for a number of cycles of laser illumination. This material is available free of charge via the Internet at http://pubs.acs.org.

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